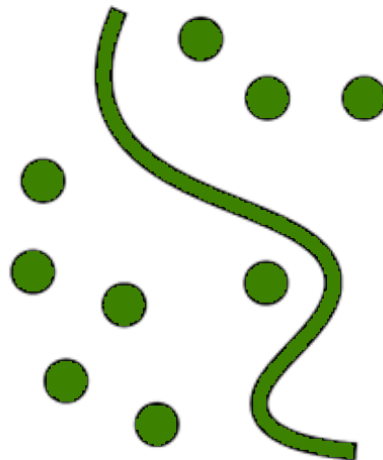


Generic event classification using rapidity-mass matrices and machine learning

<https://arxiv.org/abs/1805.11650>

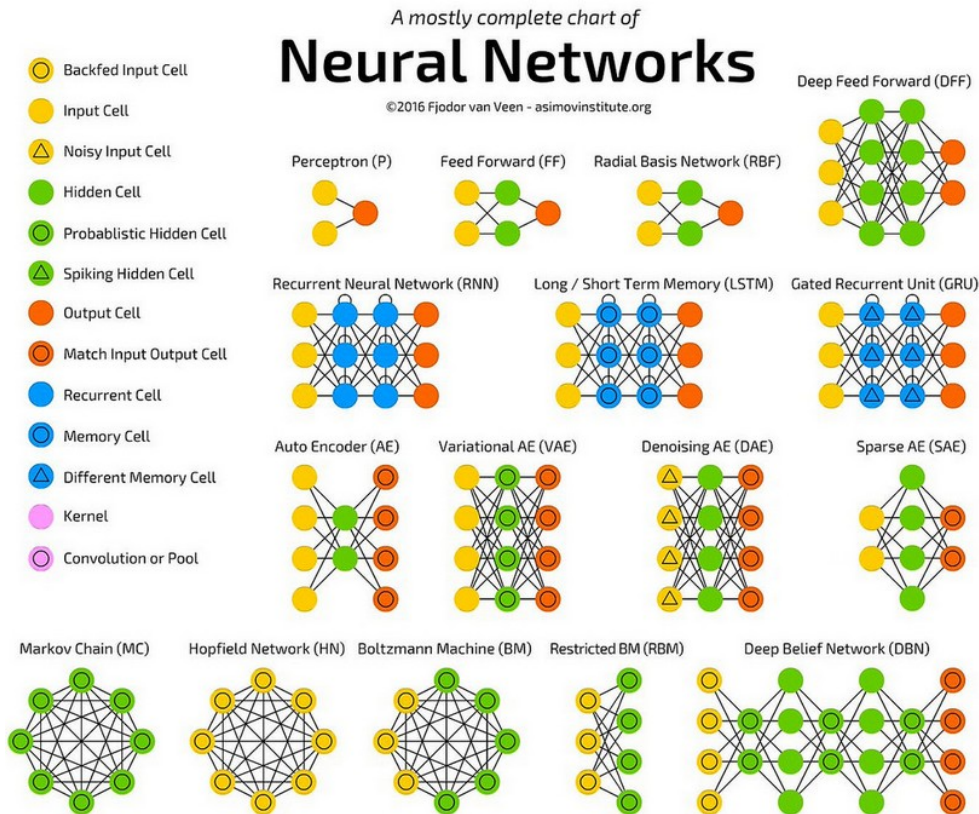
S. Chekanov

ANL Theory Seminar
(September 25, 2018)



Using Artificial Neural Networks (ANN)

- Many new neural networks invented in recent years
- Pushed by leading industries for object identification in images/video



When applies to HEP:

Success of ANN in HEP is loosely related to the choice ANN (i.e. one ANN vs another ANN)

How to prepare a meaningful input?

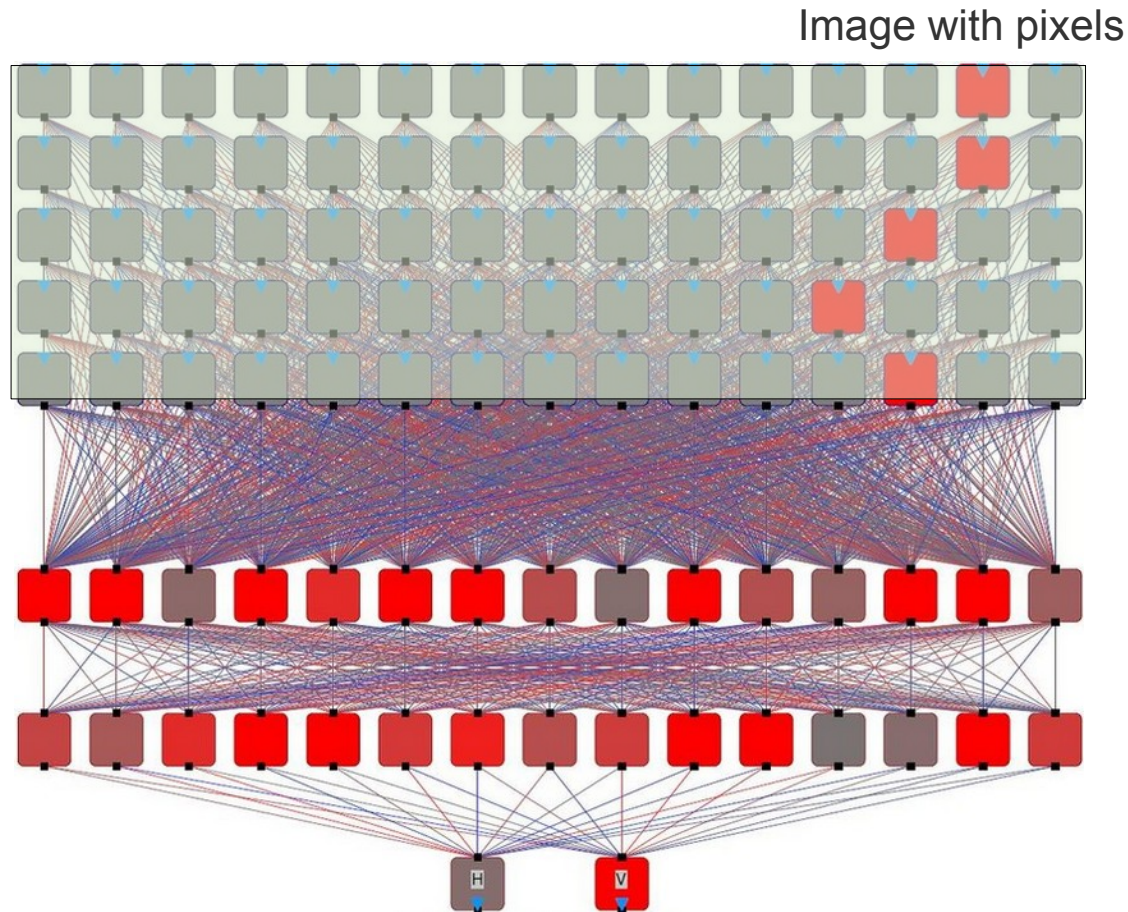
90% of time spent on hand-crafting input variables

Credit the Asimov Institute & J.Stirrup

Event classification using imaging of collision events. S.Chekanov (ANL) et al

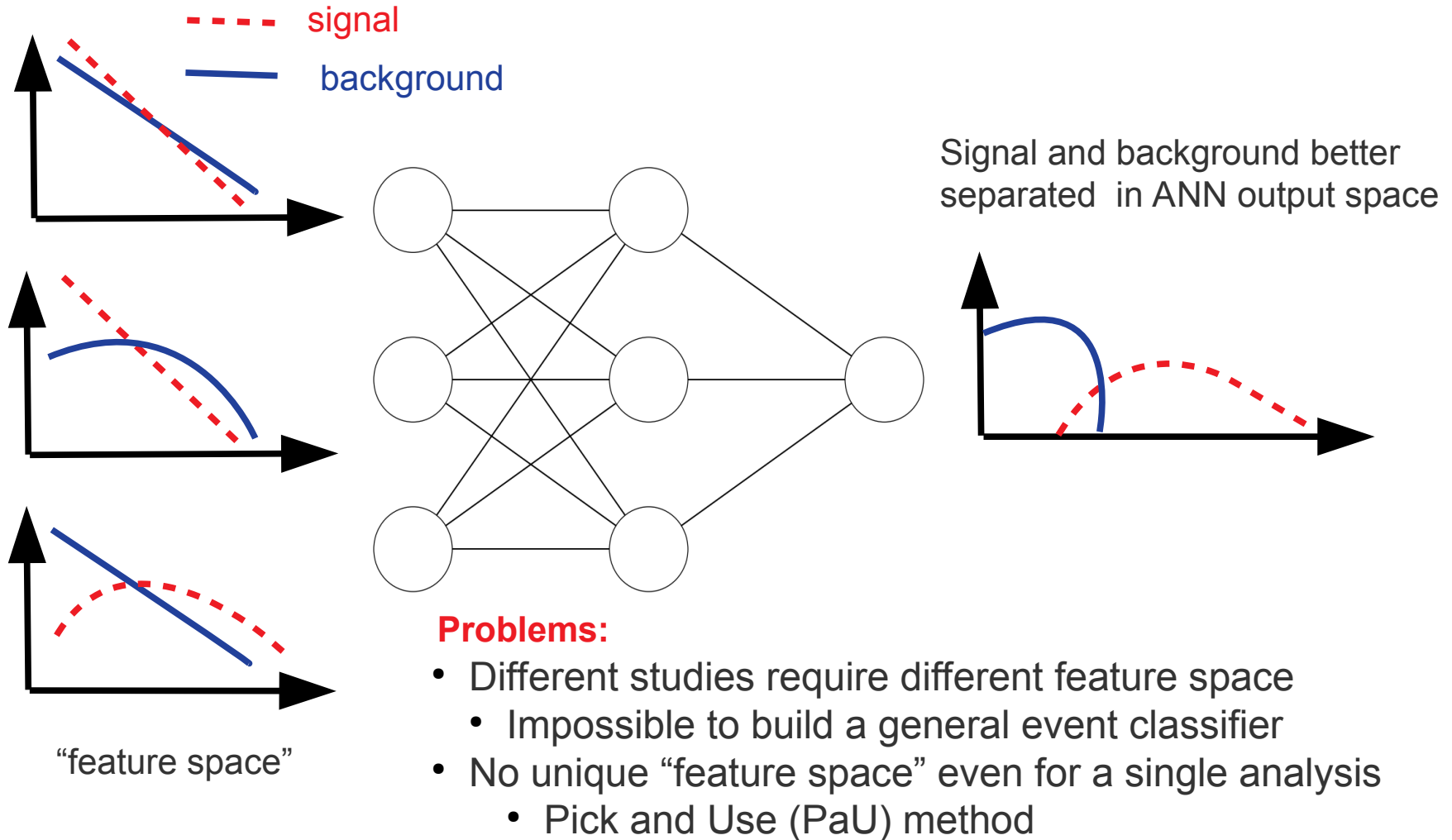


Most common usage of ANN: Image identification



ANN in HEP

- Extensively used in HEP in the last ~25 years



Feature space for event classifications

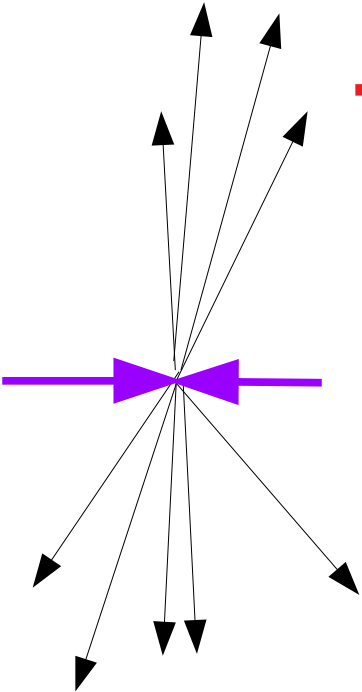
- **Event classification depends on prepared inputs**
 - Identify variables with background and signal “features”
 - Data and dimensionality reduction
 - Data re-scale (the range between 0 and 1 is a popular choice),
 - Data normalization (to avoid cases when some of input values overweight others)
 - etc.
- **ANN are suppose to simplify analysis but:**
 - Preparing analysis for NN is time consuming
 - Need to hand-pick variables, study them etc.. No uniqueness of input variables.
- **Idea: create a general image-like transformation of lists with 4-momenta to data structures that reflect most significant features of hadronic-final state**
 - General representation of collision event. Single and double- particle densities
 - Natural language for machine learning → leverage algorithms from leading industries
 - Easy to visualize for humans
 - Leverage algorithms for image identification from leading industries

FREE BONUS!



Mapping collision events to 2D arrays

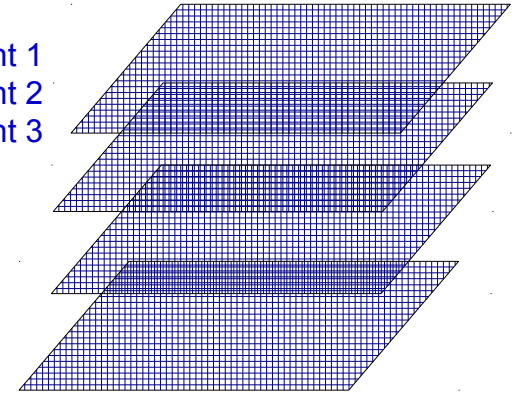
List with 4-momenta of N particles



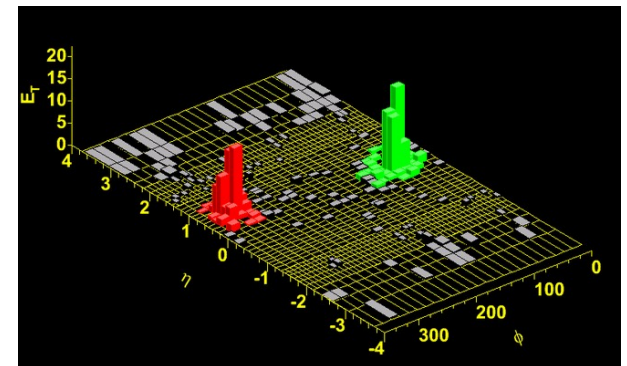
Matrices:

- Fixed size
- Dimensionless
- Lorentz invariant
- Fixed range of values
- Single particle densities
- Two-particle correlations
- Independent cells
- Cells connected by proximity due to a well-defined hierarchy
- Easy to visualize

event 1
event 2
event 3
...



NOT GOOD for our goal



Event classification using imaging of collision events. S.Chekanov (ANL) et al



Rapidity-mass matrix (RMM)

← jets →
← muons →
← .. electrons, photons →

$$\begin{pmatrix}
 e_T^{\text{miss}} & m_T(j_1) & m_T(j_2) & \dots & m_T(j_N) & m_T(\mu_1) & m_T(\mu_2) & \dots & m_T(\mu_N) \\
 h_L(j_1) & e_T(\mathbf{j}_1) & m(j_1, j_2) & \dots & m(j_1, j_N) & m(j_1, \mu_1) & m(j_1, \mu_2) & \dots & m(j_1, \mu_N) \\
 h_L(j_2) & h(j_1, j_2) & \delta e_T(\mathbf{j}_2) & \dots & m(j_2, j_N) & m(j_2, \mu_1) & m(j_2, \mu_2) & \dots & m(j_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(j_N) & h(j_1, j_N) & \dots & \dots & \delta e_T(\mathbf{j}_N) & m(j_N, \mu_1) & m(j_N, \mu_2) & \dots & m(j_N, \mu_N) \\
 h_L(\mu_1) & h(\mu_1, j_1) & h(\mu_1, j_2) & \dots & h(\mu_1, j_N) & e_T(\mu_1) & m(\mu_1, \mu_2) & \dots & m(\mu_1, \mu_N) \\
 h_L(\mu_2) & h(\mu_2, j_1) & h(\mu_2, j_2) & \dots & h(\mu_2, j_N) & h(\mu_1, \mu_2) & \delta e_T(\mu_2) & \dots & m(\mu_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(\mu_N) & h(\mu_N, j_1) & h(\mu_N, j_2) & \dots & h(\mu_N, j_N) & h(\mu_N, \mu_1) & h(\mu_N, \mu_2) & \dots & \delta e_T(\mu_N)
 \end{pmatrix}$$

e_T^{miss} – missing ET of events

$m_T(i)$ - transverse mass of object “i”

$e_T(i)$ - transverse energy (ordered)

$\delta e_T(i)$ – transverse energy imbalances

$m(i,j)$ – two-particle invariant masses

$h_L(i)$ - $\cosh(y)-1$ (y is rapidity) – Lorentz factor

$h(i,j)$ - $\cosh(0.5(y_i - y_j)) - 1$ – rapidity difference

} scaled by $1/\sqrt{s}$

} scaled by a constant

What does this matrix represent?

Events with missing particles (large missing ET)

Missing transverse energy

$$\begin{pmatrix}
 \mathbf{e}_T^{\text{miss}} & m_T(j_1) & m_T(j_2) & \dots & m_T(j_N) & m_T(\mu_1) & m_T(\mu_2) & \dots & m_T(\mu_N) \\
 h_L(j_1) & \mathbf{e}_T(\mathbf{j}_1) & m(j_1, j_2) & \dots & m(j_1, j_N) & m(j_1, \mu_1) & m(j_1, \mu_2) & \dots & m(j_1, \mu_N) \\
 h_L(j_2) & h(j_1, j_2) & \delta \mathbf{e}_T(\mathbf{j}_2) & \dots & m(j_2, j_N) & m(j_2, \mu_1) & m(j_2, \mu_2) & \dots & m(j_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(j_N) & h(j_1, j_N) & \dots & \dots & \delta \mathbf{e}_T(\mathbf{j}_N) & m(j_N, \mu_1) & m(j_N, \mu_2) & \dots & m(j_N, \mu_N) \\
 h_L(\mu_1) & h(\mu_1, j_1) & h(\mu_1, j_2) & \dots & h(\mu_1, j_N) & \mathbf{e}_T(\mu_1) & m(\mu_1, \mu_2) & \dots & m(\mu_1, \mu_N) \\
 h_L(\mu_2) & h(\mu_2, j_1) & h(\mu_2, j_2) & \dots & h(\mu_2, j_N) & h(\mu_1, \mu_2) & \delta \mathbf{e}_T(\mu_2) & \dots & m(\mu_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(\mu_N) & h(\mu_N, j_1) & h(\mu_N, j_2) & \dots & h(\mu_N, j_N) & h(\mu_N, \mu_1) & h(\mu_N, \mu_2) & \dots & \delta \mathbf{e}_T(\mu_N)
 \end{pmatrix}$$

Sensitive to particles that include decays to invisible particles

Missing transverse mass for each particle type

SM particles with neutrino decays, exotic particles (SUSY, Dark Matter candidates etc)

Example: reconstruction of transverse masses ($W \rightarrow \mu \nu$)

$$\begin{pmatrix}
 \mathbf{e}_T^{\text{miss}} & m_T(j_1) & m_T(j_2) & \dots & m_T(j_N) & m_T(\mu_1) & m_T(\mu_2) & \dots & m_T(\mu_N) \\
 h_L(j_1) & \mathbf{e}_T(\mathbf{j}_1) & m(j_1, j_2) & \dots & m(j_1, j_N) & m(j_1, \mu_1) & m(j_1, \mu_2) & \dots & m(j_1, \mu_N) \\
 h_L(j_2) & h(j_1, j_2) & \delta \mathbf{e}_T(\mathbf{j}_2) & \dots & m(j_2, j_N) & m(j_2, \mu_1) & m(j_2, \mu_2) & \dots & m(j_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(j_N) & h(j_1, j_N) & \dots & \dots & \delta \mathbf{e}_T(\mathbf{j}_N) & m(j_N, \mu_1) & m(j_N, \mu_2) & \dots & m(j_N, \mu_N) \\
 h_L(\mu_1) & h(\mu_1, j_1) & h(\mu_1, j_2) & \dots & h(\mu_1, j_N) & \mathbf{e}_T(\mu_1) & m(\mu_1, \mu_2) & \dots & m(\mu_1, \mu_N) \\
 h_L(\mu_2) & h(\mu_2, j_1) & h(\mu_2, j_2) & \dots & h(\mu_2, j_N) & h(\mu_1, \mu_2) & \delta \mathbf{e}_T(\mu_2) & \dots & m(\mu_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(\mu_N) & h(\mu_N, j_1) & h(\mu_N, j_2) & \dots & h(\mu_N, j_N) & h(\mu_N, \mu_1) & h(\mu_N, \mu_2) & \dots & \delta \mathbf{e}_T(\mu_N)
 \end{pmatrix}$$

Transverse momenta of all objects

$$\begin{pmatrix}
 \mathbf{e}_T^{\text{miss}} & m_T(j_1) & m_T(j_2) & \dots & m_T(j_N) & m_T(\mu_1) & m_T(\mu_2) & \dots & m_T(\mu_N) \\
 h_L(j_1) & \mathbf{e}_T(\mathbf{j}_1) & m(j_1, j_2) & \dots & m(j_1, j_N) & m(j_1, \mu_1) & m(j_1, \mu_2) & \dots & m(j_1, \mu_N) \\
 h_L(j_2) & h(j_1, j_2) & \delta \mathbf{e}_T(\mathbf{j}_2) & \dots & m(j_2, j_N) & m(j_2, \mu_1) & m(j_2, \mu_2) & \dots & m(j_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(j_N) & h(j_1, j_N) & \dots & \dots & \delta \mathbf{e}_T(\mathbf{j}_N) & m(j_N, \mu_1) & m(j_N, \mu_2) & \dots & m(j_N, \mu_N) \\
 h_L(\mu_1) & h(\mu_1, j_1) & h(\mu_1, j_2) & \dots & h(\mu_1, j_N) & \mathbf{e}_T(\mu_1) & m(\mu_1, \mu_2) & \dots & m(\mu_1, \mu_N) \\
 h_L(\mu_2) & h(\mu_2, j_1) & h(\mu_2, j_2) & \dots & h(\mu_2, j_N) & h(\mu_1, \mu_2) & \delta \mathbf{e}_T(\mu_2) & \dots & m(\mu_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(\mu_N) & h(\mu_N, j_1) & h(\mu_N, j_2) & \dots & h(\mu_N, j_N) & h(\mu_N, \mu_1) & h(\mu_N, \mu_2) & \dots & \delta \mathbf{e}_T(\mu_N)
 \end{pmatrix}$$

Diagonal elements:

Transverse momentum of leading in E_T objects

Transverse momentum imbalances

$$\delta e_T(i_n) = \frac{E_T(i_{n-1}) - E_T(i_n)}{E_T(i_{n-1}) + E_T(i_n)}$$

Transverse energy imbalances:

- Sensitive to interactions of partons in the medium of heavy ion collisions
- Can be used to separate dijet QCD events from more complex events

Can be use reconstruct:

- Transverse energies of all objects
- H_T of events
- Energies $E(i) = e_T(i) \sqrt{s} (h_L(i)/C + 1)$

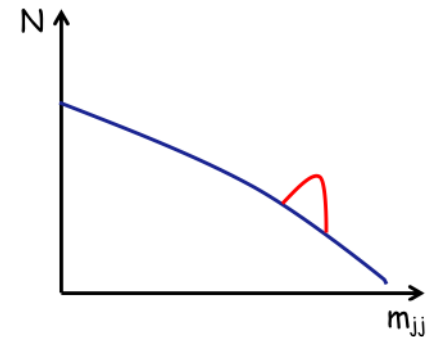
Invariant masses of particle (jet) pairs

$$\begin{pmatrix}
 \mathbf{e}_T^{\text{miss}} & m_T(j_1) & m_T(j_2) & \dots & m_T(j_N) & m_T(\mu_1) & m_T(\mu_2) & \dots & m_T(\mu_N) \\
 h_L(j_1) & \mathbf{e}_T(\mathbf{j}_1) & m(j_1, j_2) & \dots & m(j_1, j_N) & m(j_1, \mu_1) & m(j_1, \mu_2) & \dots & m(j_1, \mu_N) \\
 h_L(j_2) & h(j_1, j_2) & \delta \mathbf{e}_T(\mathbf{j}_2) & \dots & m(j_2, j_N) & m(j_2, \mu_1) & m(j_2, \mu_2) & \dots & m(j_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(j_N) & h(j_1, j_N) & \dots & \dots & \delta \mathbf{e}_T(\mathbf{j}_N) & m(j_N, \mu_1) & m(j_N, \mu_2) & \dots & m(j_N, \mu_N) \\
 h_L(\mu_1) & h(\mu_1, j_1) & h(\mu_1, j_2) & \dots & h(\mu_1, j_N) & \mathbf{e}_T(\mu_1) & m(\mu_1, \mu_2) & \dots & m(\mu_1, \mu_N) \\
 h_L(\mu_2) & h(\mu_2, j_1) & h(\mu_2, j_2) & \dots & h(\mu_2, j_N) & h(\mu_1, \mu_2) & \delta \mathbf{e}_T(\mu_2) & \dots & m(\mu_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(\mu_N) & h(\mu_N, j_1) & h(\mu_N, j_2) & \dots & h(\mu_N, j_N) & h(\mu_N, \mu_1) & h(\mu_N, \mu_2) & \dots & \delta \mathbf{e}_T(\mu_N)
 \end{pmatrix}$$

$\mathbf{m}(i,j)$ - Invariant masses of all objects (jets and identified particles)

Peaks in invariant masses will be shown as enhanced top-right cells

“Gold- standard” for search for new physics



Longitudinal flow in events

$$\begin{pmatrix}
 \mathbf{e}_T^{\text{miss}} & m_T(j_1) & m_T(j_2) & \dots & m_T(j_N) & m_T(\mu_1) & m_T(\mu_2) & \dots & m_T(\mu_N) \\
 h_L(j_1) & \mathbf{e}_T(\mathbf{j}_1) & m(j_1, j_2) & \dots & m(j_1, j_N) & m(j_1, \mu_1) & m(j_1, \mu_2) & \dots & m(j_1, \mu_N) \\
 h_L(j_2) & h(j_1, j_2) & \delta \mathbf{e}_T(\mathbf{j}_2) & \dots & m(j_2, j_N) & m(j_2, \mu_1) & m(j_2, \mu_2) & \dots & m(j_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(j_N) & h(j_1, j_N) & \dots & \dots & \delta \mathbf{e}_T(\mathbf{j}_N) & m(j_N, \mu_1) & m(j_N, \mu_2) & \dots & m(j_N, \mu_N) \\
 h_L(\mu_1) & h(\mu_1, j_1) & h(\mu_1, j_2) & \dots & h(\mu_1, j_N) & \mathbf{e}_T(\mu_1) & m(\mu_1, \mu_2) & \dots & m(\mu_1, \mu_N) \\
 h_L(\mu_2) & h(\mu_2, j_1) & h(\mu_2, j_2) & \dots & h(\mu_2, j_N) & h(\mu_1, \mu_2) & \delta \mathbf{e}_T(\mu_2) & \dots & m(\mu_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(\mu_N) & h(\mu_N, j_1) & h(\mu_N, j_2) & \dots & h(\mu_N, j_N) & h(\mu_N, \mu_1) & h(\mu_N, \mu_2) & \dots & \delta \mathbf{e}_T(\mu_N)
 \end{pmatrix}$$

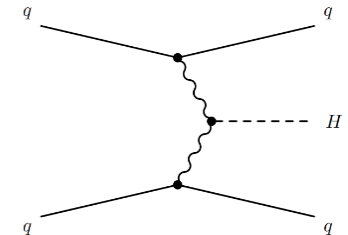
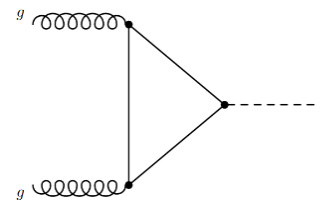
$$h_L = \cosh(y) - 1:$$

- Large values for forward physics
- Small value (~ 0) for central production

$$\cosh(y) = \gamma \text{ is Lorentz factor}$$

Example:

- VBF Higgs production has large values in the first column (centrally produced)
- ..



Rapidity differences

$$\begin{pmatrix}
 \mathbf{e}_T^{\text{miss}} & m_T(j_1) & m_T(j_2) & \dots & m_T(j_N) & m_T(\mu_1) & m_T(\mu_2) & \dots & m_T(\mu_N) \\
 h_L(j_1) & \mathbf{e}_T(\mathbf{j}_1) & m(j_1, j_2) & \dots & m(j_1, j_N) & m(j_1, \mu_1) & m(j_1, \mu_2) & \dots & m(j_1, \mu_N) \\
 h_L(j_2) & h(j_1, j_2) & \delta \mathbf{e}_T(\mathbf{j}_2) & \dots & m(j_2, j_N) & m(j_2, \mu_1) & m(j_2, \mu_2) & \dots & m(j_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(j_N) & h(j_1, j_N) & \dots & \dots & \delta \mathbf{e}_T(\mathbf{j}_N) & m(j_N, \mu_1) & m(j_N, \mu_2) & \dots & m(j_N, \mu_N) \\
 h_L(\mu_1) & h(\mu_1, j_1) & h(\mu_1, j_2) & \dots & h(\mu_1, j_N) & \mathbf{e}_T(\mu_1) & m(\mu_1, \mu_2) & \dots & m(\mu_1, \mu_N) \\
 h_L(\mu_2) & h(\mu_2, j_1) & h(\mu_2, j_2) & \dots & h(\mu_2, j_N) & h(\mu_1, \mu_2) & \delta \mathbf{e}_T(\mu_2) & \dots & m(\mu_2, \mu_N) \\
 \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots & \dots \\
 h_L(\mu_N) & h(\mu_N, j_1) & h(\mu_N, j_2) & \dots & h(\mu_N, j_N) & h(\mu_N, \mu_1) & h(\mu_N, \mu_2) & \dots & \delta \mathbf{e}_T(\mu_N)
 \end{pmatrix}$$

$$h(i,j) \sim \cosh (y_j - y_i) - 1$$

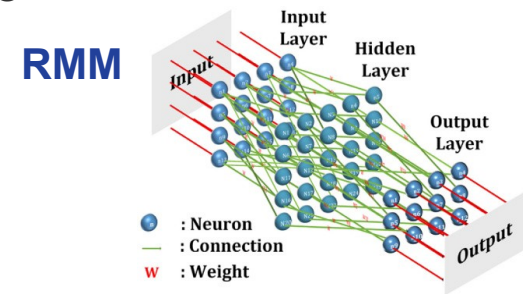
$$h(i,j) \sim 0 \text{ for collimated 2-particles}$$

Rapidity difference used for:

- Dijet searches (CMS, ATLAS)
- Probing parton dynamics (CMS)

Useful features

- Dimensionless, Lorentz invariant (except for Lorentz factors for all objects)
- Small linear correlations between RMM cells
 - No redundant information
 - According to Monte Carlo simulations:
 - Pearson correlation coefficient >0.5 for 0.3% of cell pairs
 - Out of 0.4%, 50% correlation is seen between m_T and $m(i,j)$
- Well-defined hierarchy by construction
 - Cells connected by proximity
 - Should look as “images”, not as random noise.. → good for visualization
- Natural language for machine learning:
 - Each cell maps to a fixed node/neuron.
 - Normalization and standardization



However:


- *RMM is a sparse matrix for single events. Keep non-zero values and their indexes!*
- *Averaging over many events makes more visually appealing images*

Monte Carlo simulations

Several processes from Pythia8 (LO+PS)

- **Dijet QCD:**
 - All $2 \rightarrow 2$ processes (10)
- **Top production:**
 - $g g \rightarrow t \bar{t}$
 - $q \bar{q} \rightarrow t \bar{t}$
- **Charged Higgs production**
 - $b g \rightarrow H^\pm t$
- **Double boson production**
 - $f \bar{f} \rightarrow \gamma^*/Z \gamma^*/Z$
 - $f \bar{f}' \rightarrow Z W^\pm$
 - $f \bar{f} \rightarrow W^+ W^-$
- **SM Higgs production**

<http://atlaswww.hep.anl.gov/hepsim/>



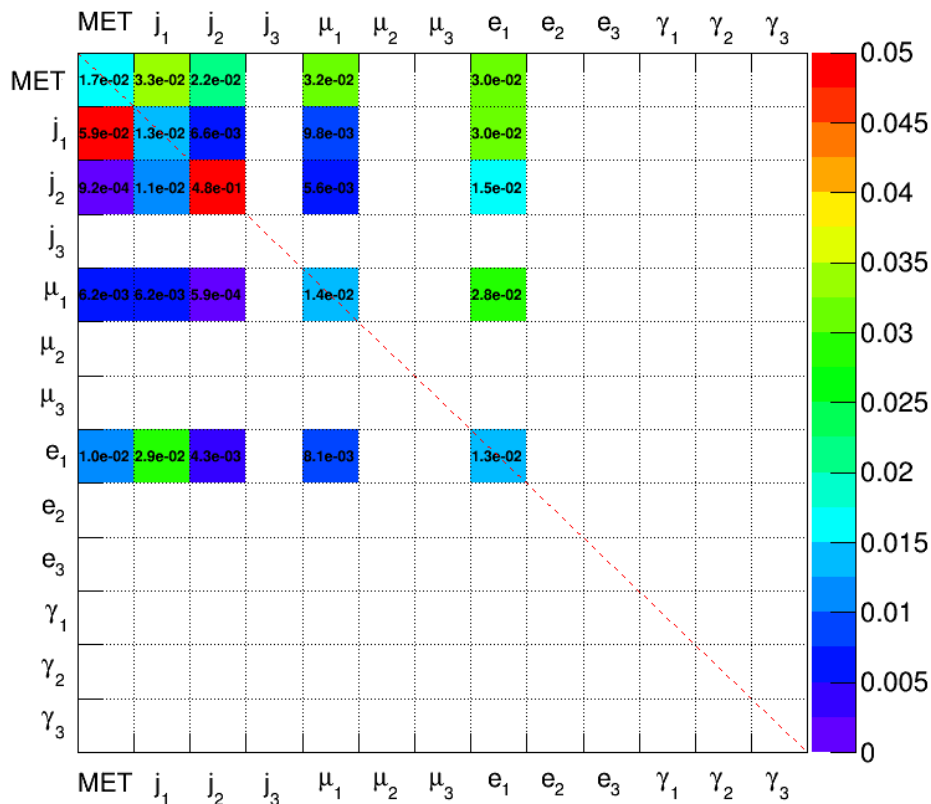
The screenshot shows the HepSim website interface. At the top, there is a navigation bar with links: "Get involved", "Full Search", "Experiments", "Manual", "Mirrors", "Tools", "About", and "Login". The main heading is "HepSim" with the subtitle "Repository with Monte Carlo simulations for particle physics". A search bar is visible on the right. Below the heading, there is a table listing simulation datasets. The table has columns: "Id", "E [TeV]", "Dataset name", "Generator", "Process", "Topic", "Files", and "Created". The table contains 13 rows of data, including processes like "Various SM/BSM process for ML", "QCD dijets for ProIO tests", "QCD dijets for tests of ProIO", "DIS events at Q2>1 GeV2", "H- top with H- to HW and tan(beta)=1-7", "H- top with H- to tb and tan(beta)=1-7", "H+ W- with H+ decay to t-bbar tan(beta)=1-7", "H+ W- with H+ decay to HW for tan(beta)=1-7", and "Higgs to gamma gamma".

Id	E [TeV]	Dataset name	Generator	Process	Topic	Files	Created
328	pp 13	tev13pp_pythia8_rmm	PYTHIA8	Various SM/BSM process for ML	SM	Info	2018/09/16
327	pp 13	tev13pp_qcd_pythia8_proio	PYTHIA8	QCD dijets for ProIO tests	SM	Info	2018/08/27
326	pp 13	tev13pp_qcd_pythia8_proio_tests	PYTHIA8	QCD dijets for tests of ProIO	SM	Info	2018/08/20
325	e-p 0.035	gev35ep_pythia8_dis1q2ct14lo	PYTHIA8	DIS events at Q2>1 GeV2	SM	Info	2018/07/25
323	pp 13	tev13pp_mg5_chaHT_tbeta_hw	MADGRAPH/PY8	H- top with H- to HW and tan(beta)=1-7	Exotics	Info	2018/06/13
322	pp 13	tev13pp_mg5_chaHT_tbeta_tb	MADGRAPH/PY8	H- top with H- to tb and tan(beta)=1-7	Exotics	Info	2018/06/13
321	pp 13	tev13pp_mg5_chaHW_tbeta_tb	MADGRAPH/PY8	H+ W- with H+ decay to t-bbar tan(beta)=1-7	Exotics	Info	2018/06/06
320	pp 13	tev13pp_mg5_chaHW_tbeta_hw	MADGRAPH/PY8	H+ W- with H+ decay to HW for tan(beta)=1-7	Exotics	Info	2018/06/06
318	pp 13	tev13pp_pythia8_gamgam	PYTHIA8	Higgs to gamma gamma	SM	Info	2018/04/20

All LO processes and all top/W/H decays enabled

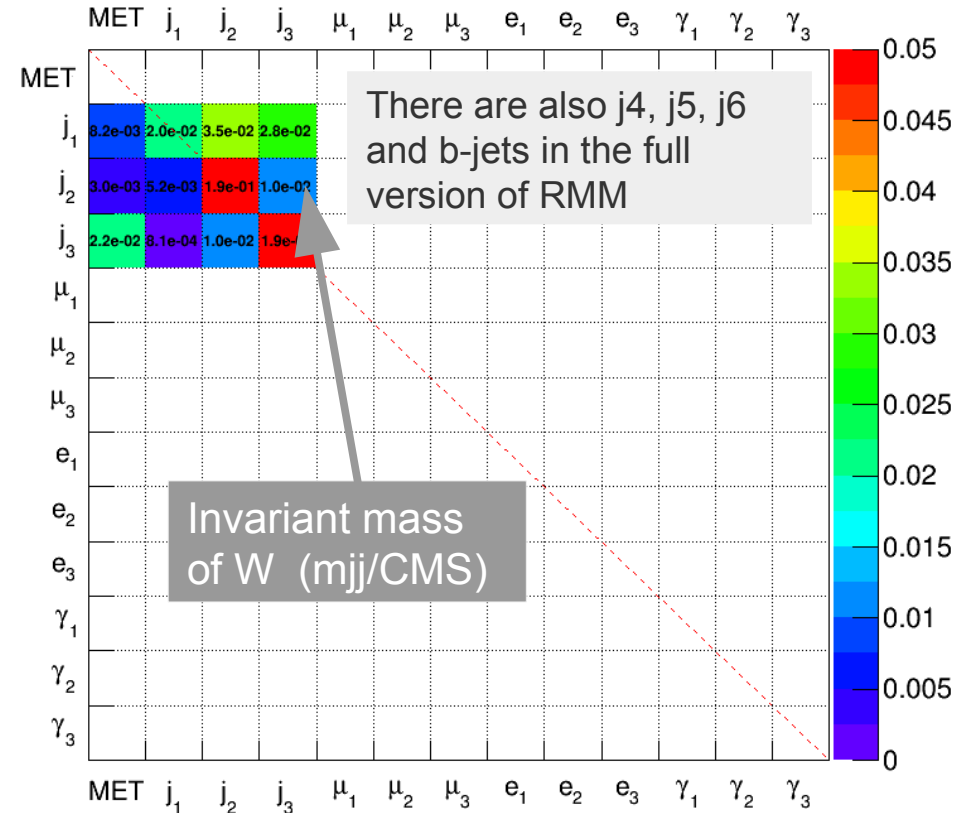
Example: Two events with $t\bar{t}$

$t\bar{t} \rightarrow Wb W\bar{b} \rightarrow e \nu b \mu \nu \bar{b}$



Cell with MET is “fired”.
Also μ and e leptons

$t\bar{t} \rightarrow Wb W\bar{b} \rightarrow 6 \text{ jets}$

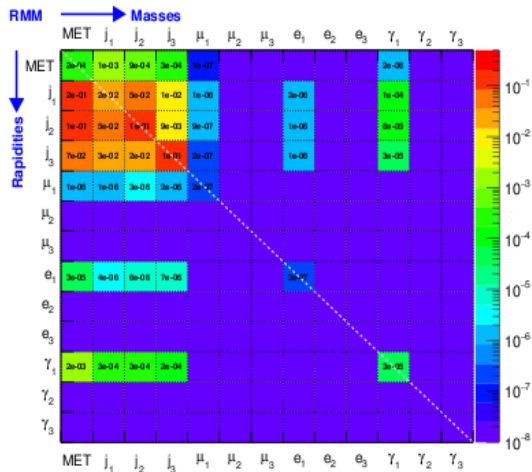


There are also j_4, j_5, j_6 and b -jets in the full version of RMM

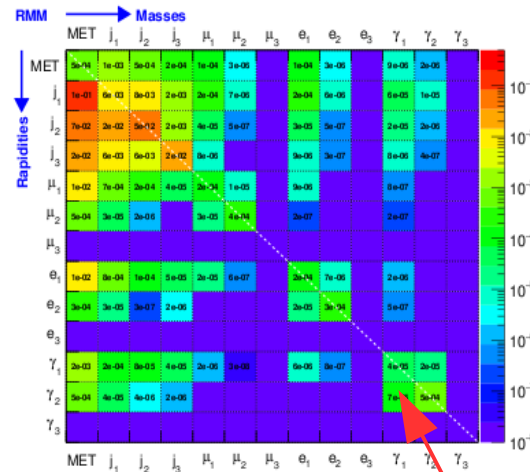
Invariant mass of W (m_{jj}/CMS)

No MET and leptons
But many jets

Average values of cells for 4 processes



(a) multijets QCD



(b) Higgs processes

Considered:

- jets, muons, electrons, photons
- up to 3 objects

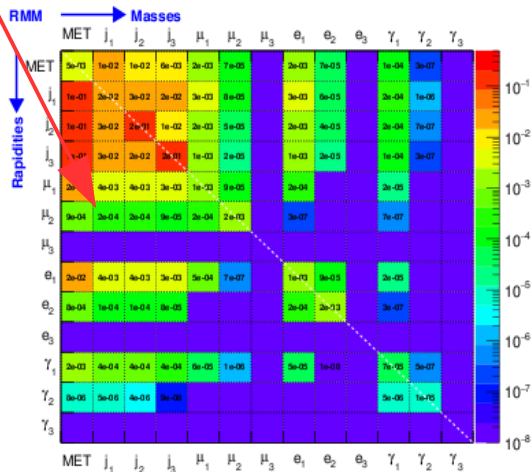
Pythia8 simulation:

All QCD processes included with all decays of top, W, H
50,000 events per RMM

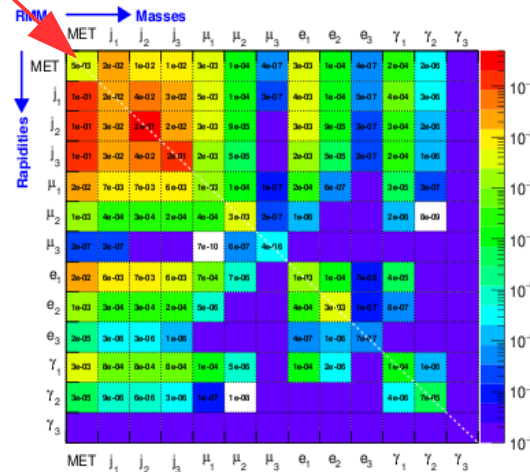
Muons

large MET

Higgs mass ($\gamma\gamma$)



(c) Top production



(d) H^+t production

$H^+ t$ is similar to $t\bar{t}$

This similarity was made intentionally by allowing $H^+ \rightarrow W H$, where $H \rightarrow b\bar{b}$

Practical example: Separation of H^+ from $t\bar{t}$

Reconstruct invariant mass of 2 jets for $t\bar{t}$ (background) and H^+ (signal)

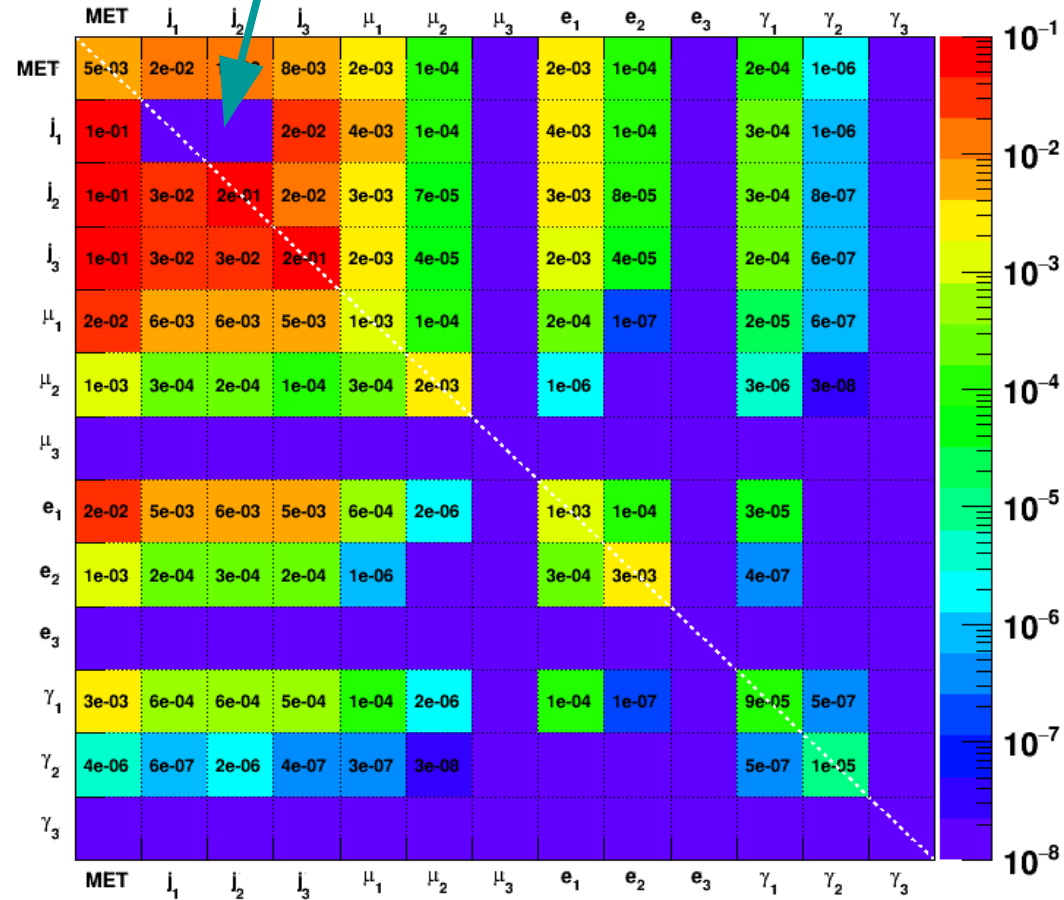
H^+ (600 GeV) \rightarrow $H W$
where $H \rightarrow b\bar{b}$

Should see a bump at ~ 600 GeV for H^+ events

Invariant mass can be calculated from RMM itself by summing up cells at (3,2) for all RMM (and scale by 13000)

To avoid biases, disable cells (2,2) and (3,2) during the NN training

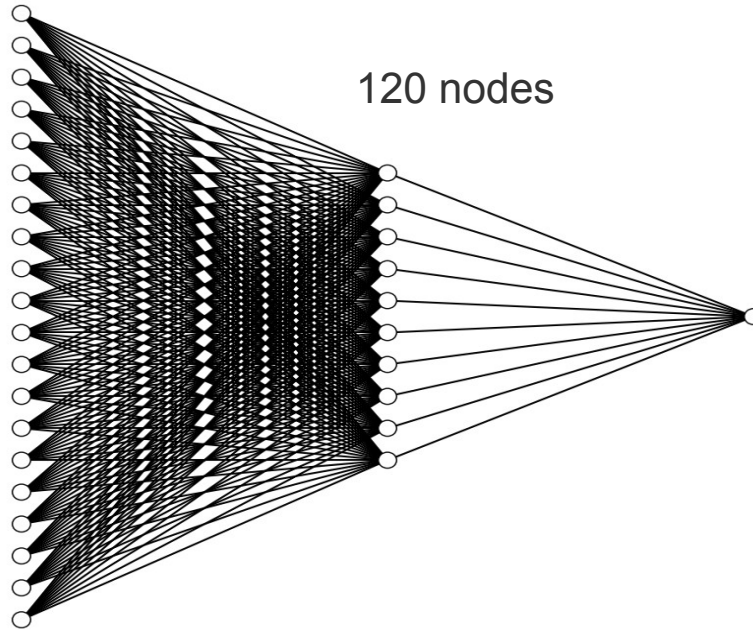
disabled links during the NN training since they are extracted "features"



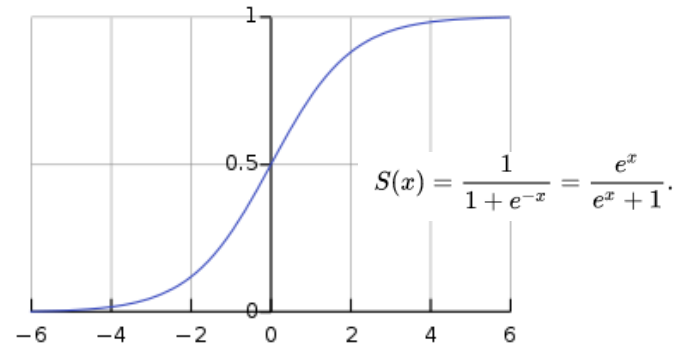
Using RMM for ANN

A simple backpropagation NN with “sigmoid” activation

169 nodes



120 nodes



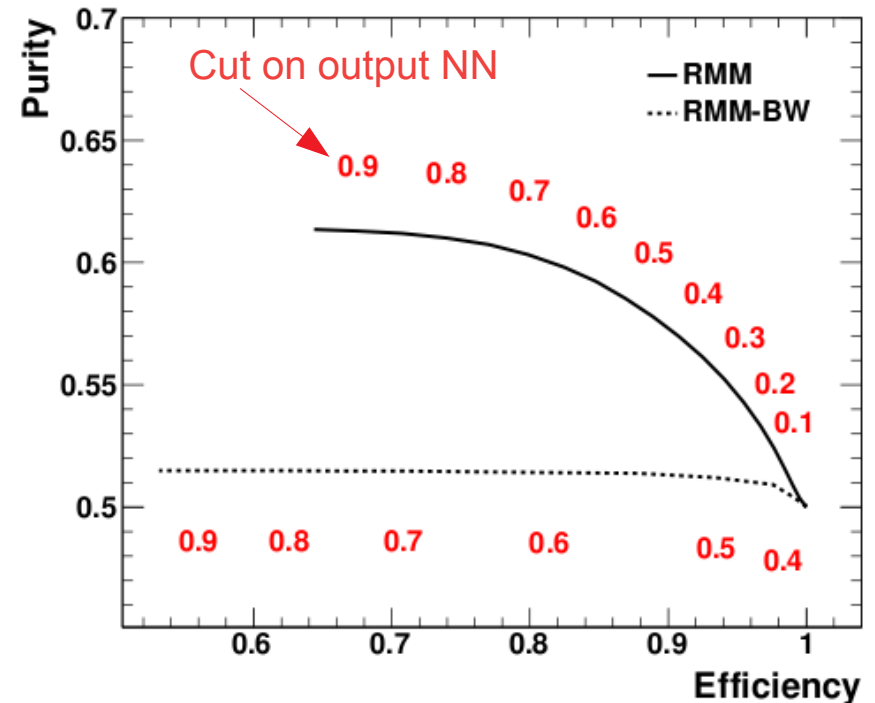
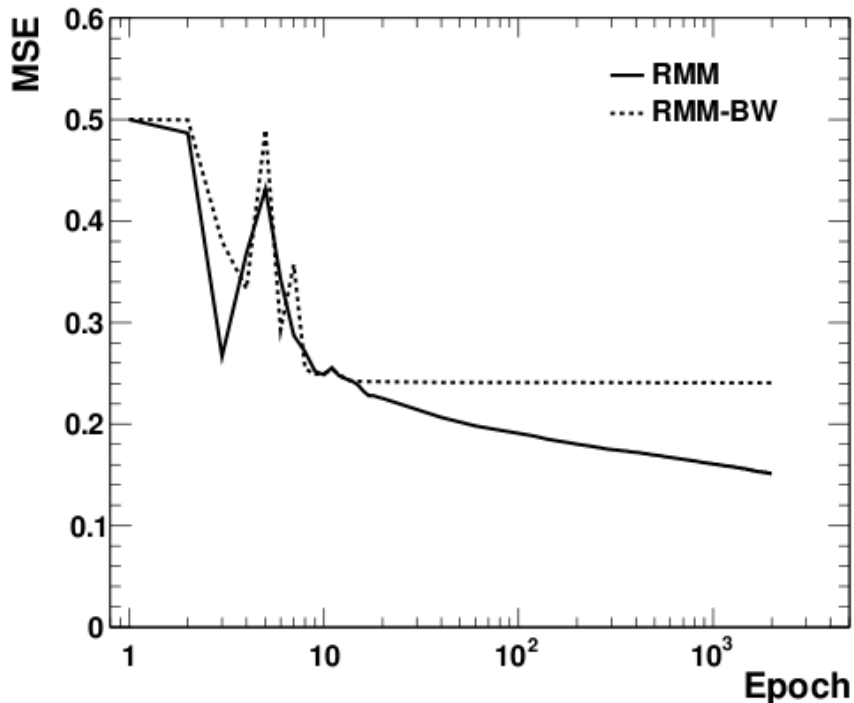
output: 0 ($t\bar{t}$) or 1 (H+)

10k Pythia8 events used to create 10k RMM (13x13)

- Use 10k events with $t\bar{t}$, and 10k with H+
- Create cross validation for ANN
- Stop training when MSE < than for cross validated ANN



Using RMM for neural networks

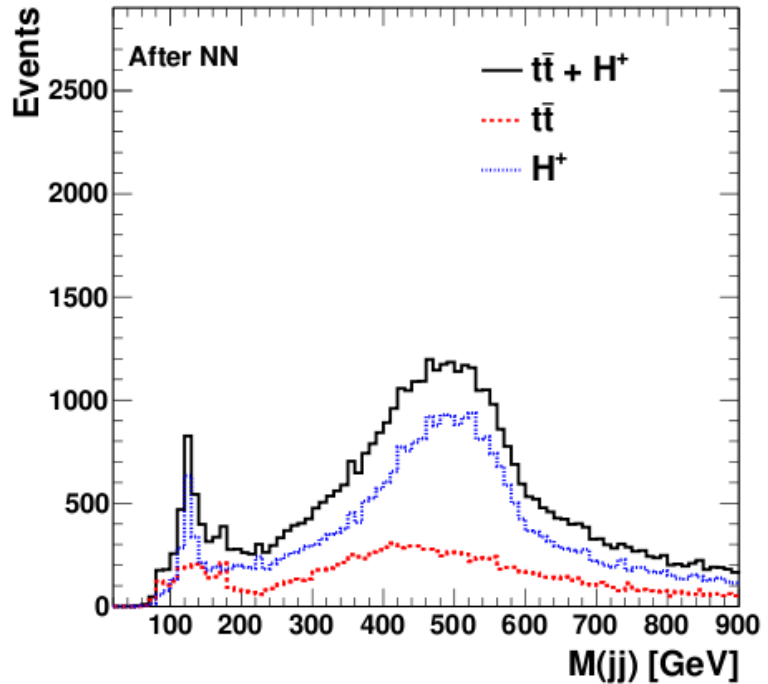
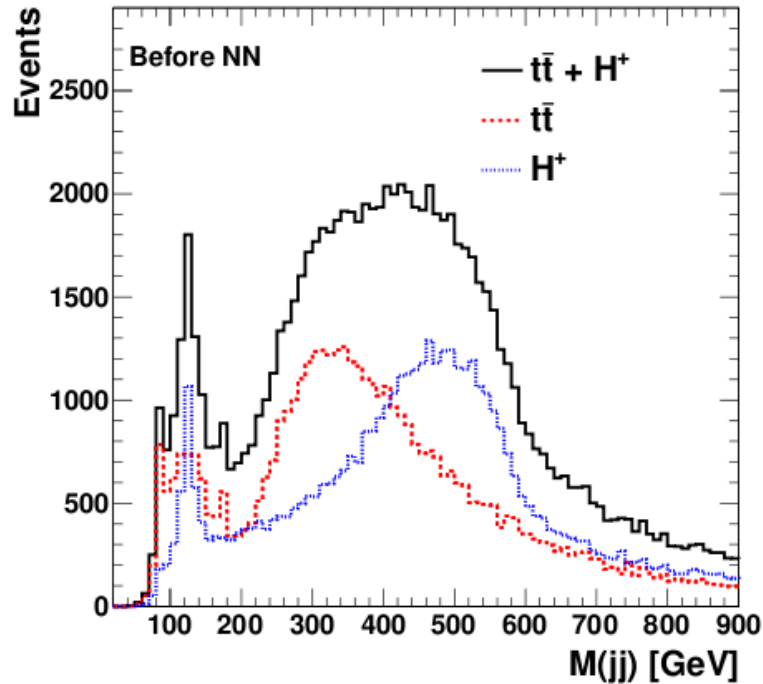


Well trained:

Mean Squared Error (MSE) decreased even in the case when only “activated” cells are used, i.e. without the actual values (dashed)
→ analogy with “black-and-white” images (RMM-BW)



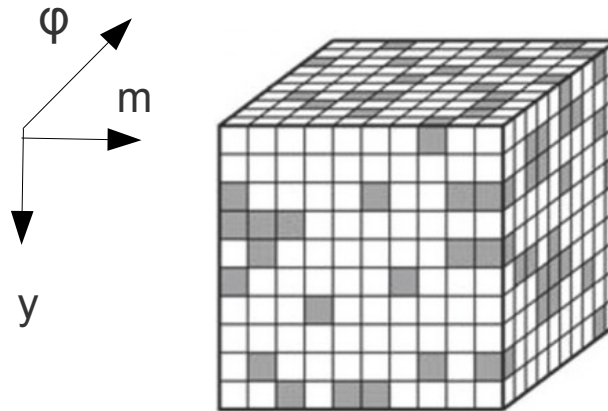
Practical example: Separation of H^+ from $t\bar{t}$



- The NN based on RMM helps to reduce S/B by a factor 3.
- Signal efficiency is reduced by 30%
- Small shift for $t\bar{t}$ (may require better tuning of disabled links)

Extending RMM

- RMM includes information on single and two-particle densities
 - but no phi due to rotational symmetry)
- Can be extended to 3D matrices to include φ , 3-particle densities etc.



Plus:

- Add tau, leptons with + and - charges (separately), b-jets
- Increase multiplicity of each object to $\sim 10-20$ (empty cells are not stored)
- Add more complex (and well reconstructed) types: J/Phi, W, Z, Higgs

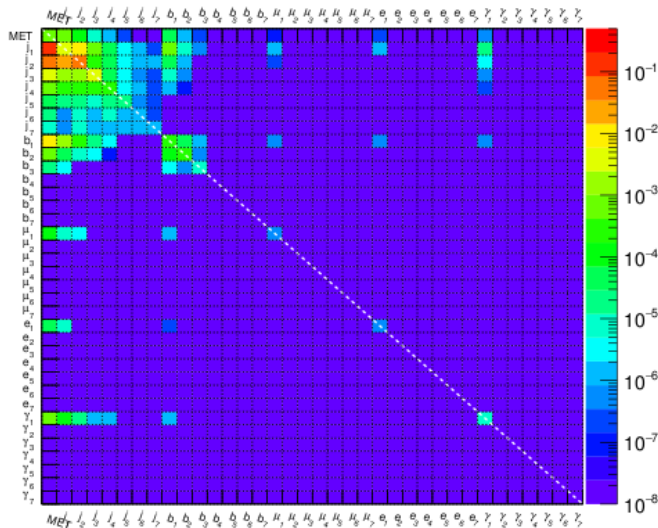


Generic event identification

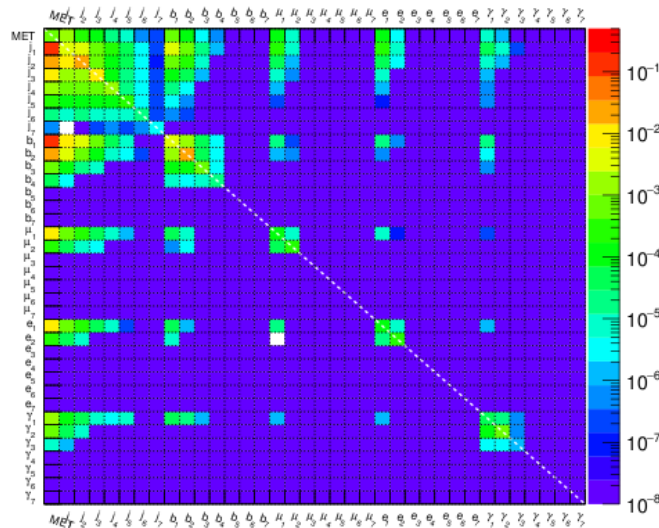
- Premise of the RMM - generality. Includes single & 2-particle densities
- No need to hand-pick input variables for every event topology/decay
- Good choice for general event classifiers?

Example:

- 5 processes with (1) SM QCD (2) Higgs (3) H+ (4) ttbar (5) Double bosons
- Create RMM using $N_p=7$ and 6 objects using b-jets



Multi-jet QCD

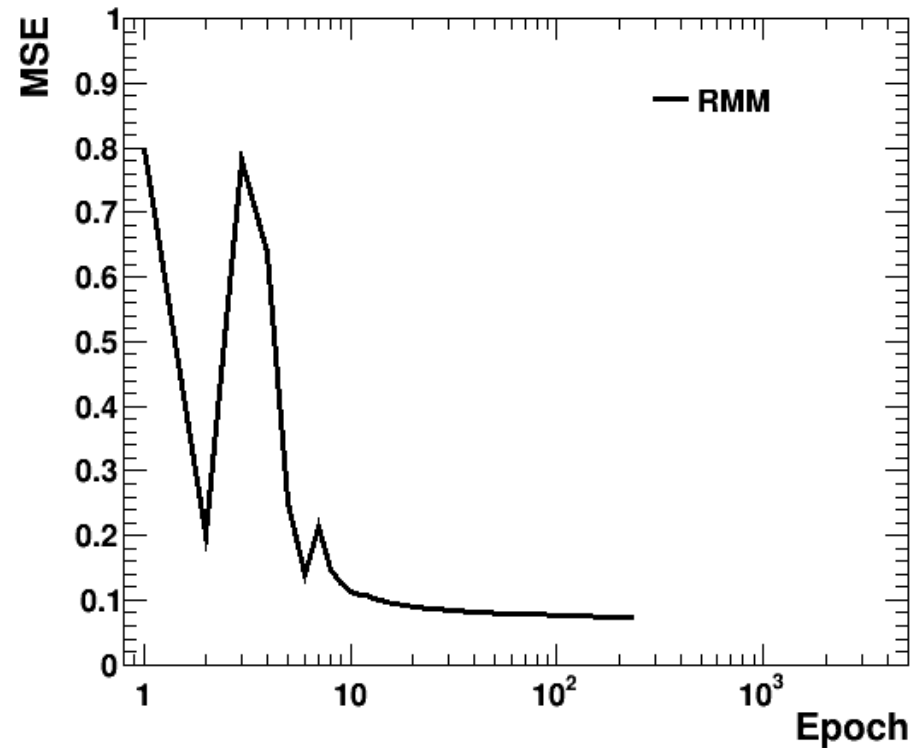
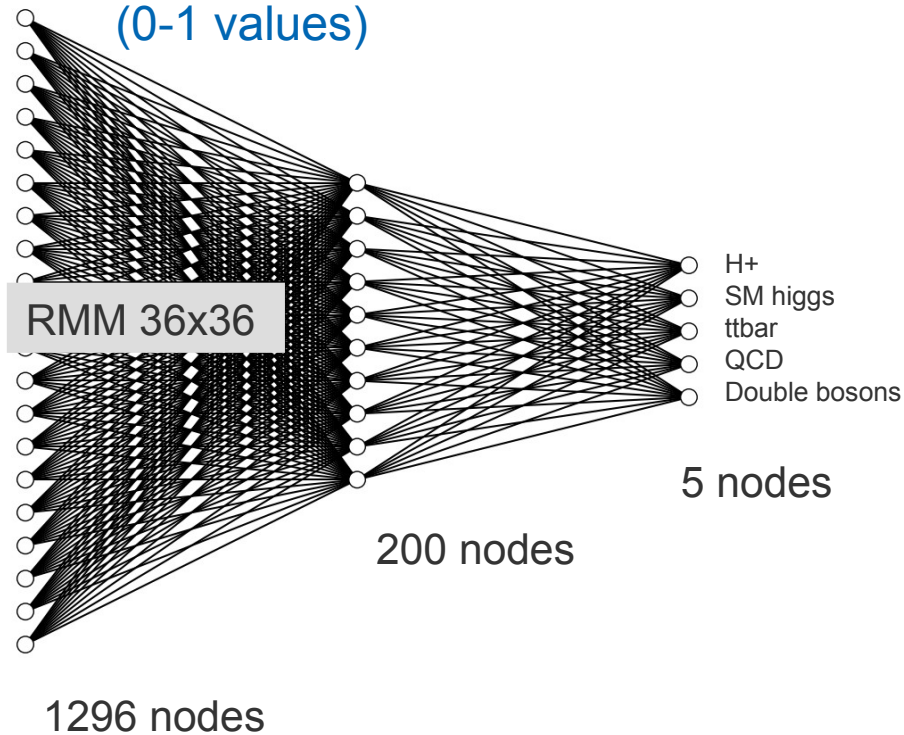


Higgs productions (all decays)

Shows average values of cells for 50k events

ANN training

Backpropagation NN with Sigmoid function, 5 outputs for each process (0-1 values)

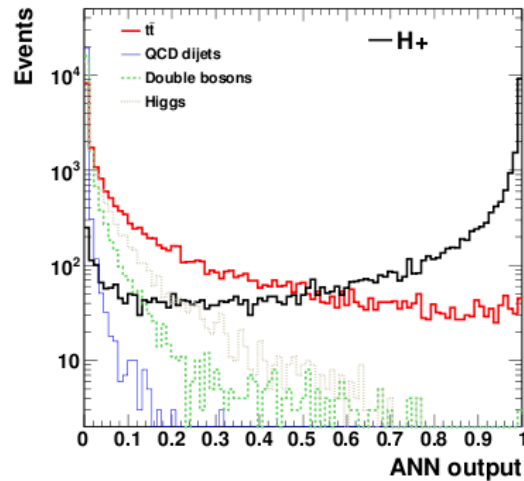


Well trained:
Mean Squared Error (MSE)
decreased from 0.8 to 0.07

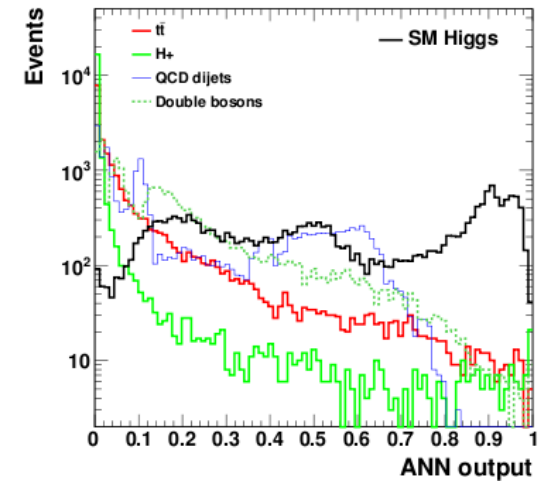


Result of NN training

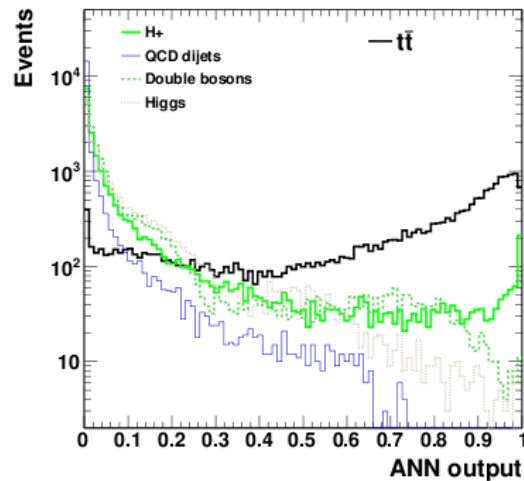
Good event separation of “signal” events (black line) from other processes



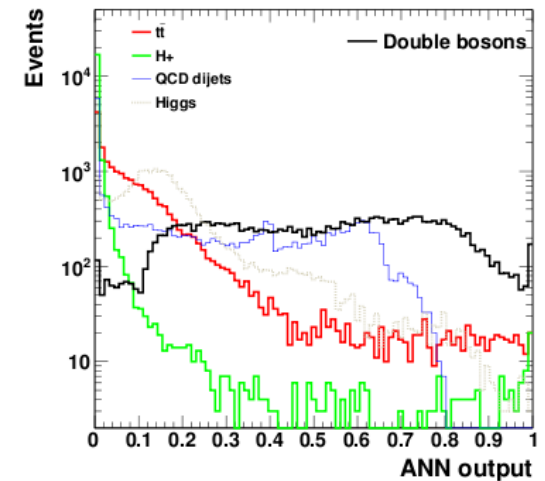
(a) Charged H^+



(b) SM Higgs



(c) $t\bar{t}$ production

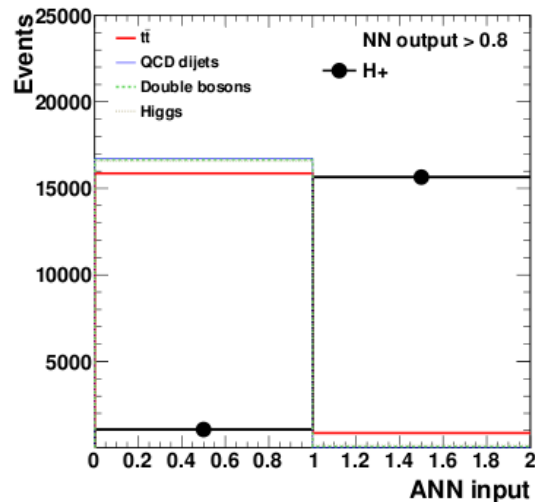


(d) Double W/Z production

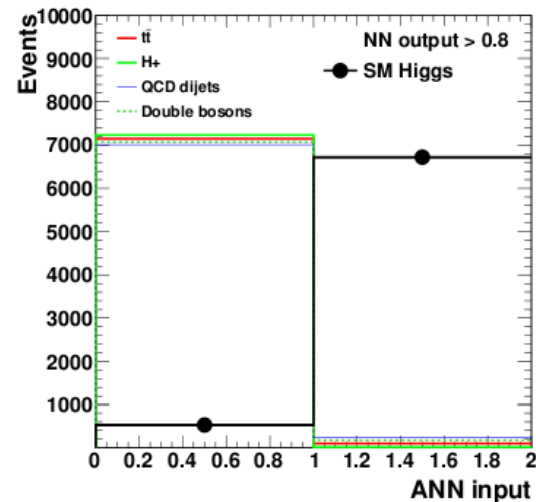


Purity of event identification

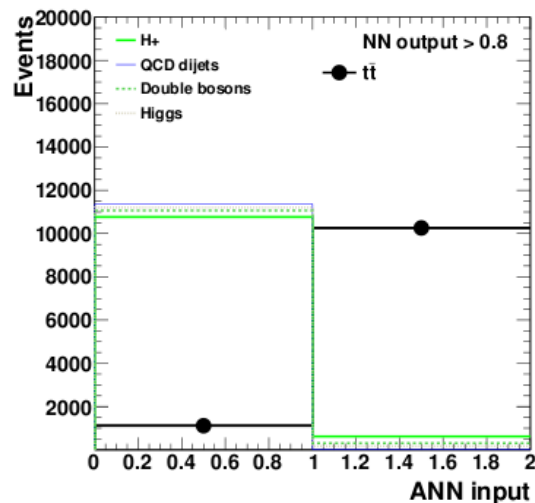
Purity of event classification is 80%-90% assuming 0.5 cut on output node



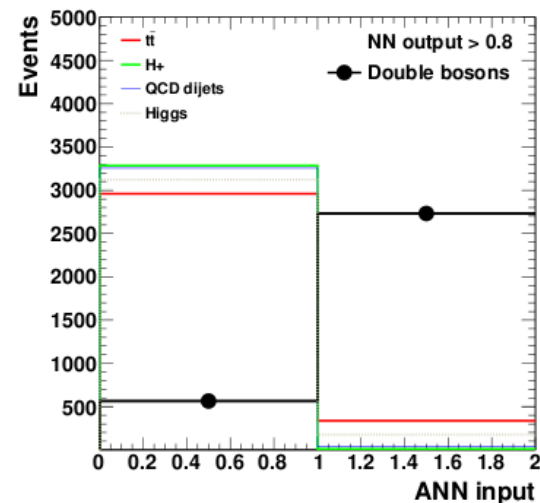
(a) Charged H⁺



(b) SM Higgs



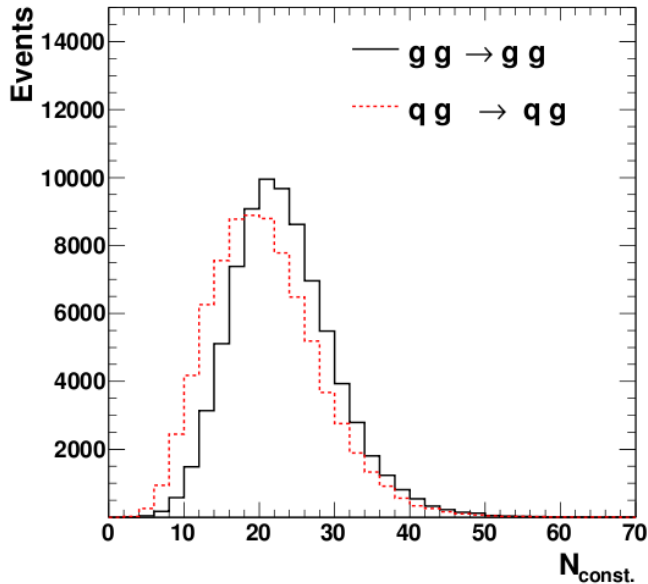
(c) $t\bar{t}$ production



(d) Double W/Z production

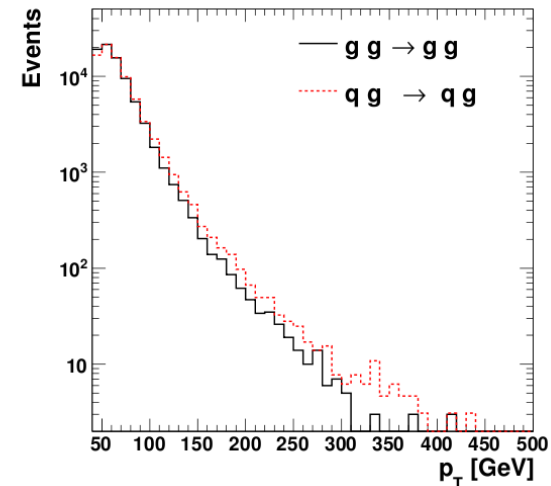
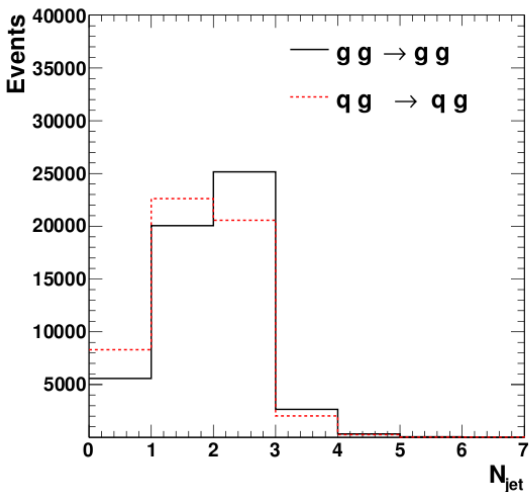
QCD dijet separation: Challenging case!

Separate gg from qg final states (dijets) \rightarrow Distributions are nearly identical. Presence of g instead of q leads to broader jets and changes in jet kinematics / shape



Well-known difference: Number of jet constituents is larger for gluon jets than for quark jets due to difference in color factors ($C_A = 3$ vs $C_F = 3/4$)

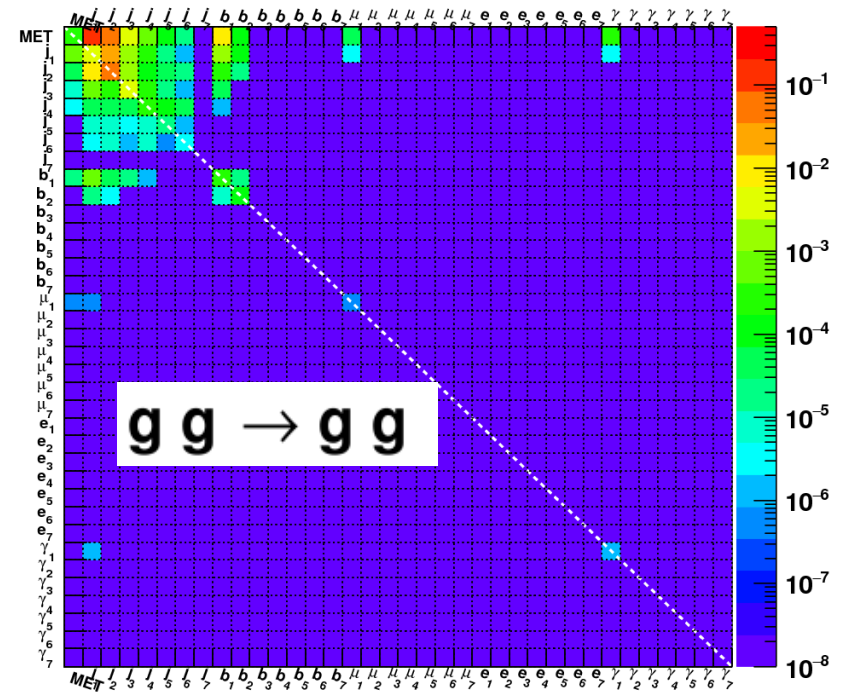
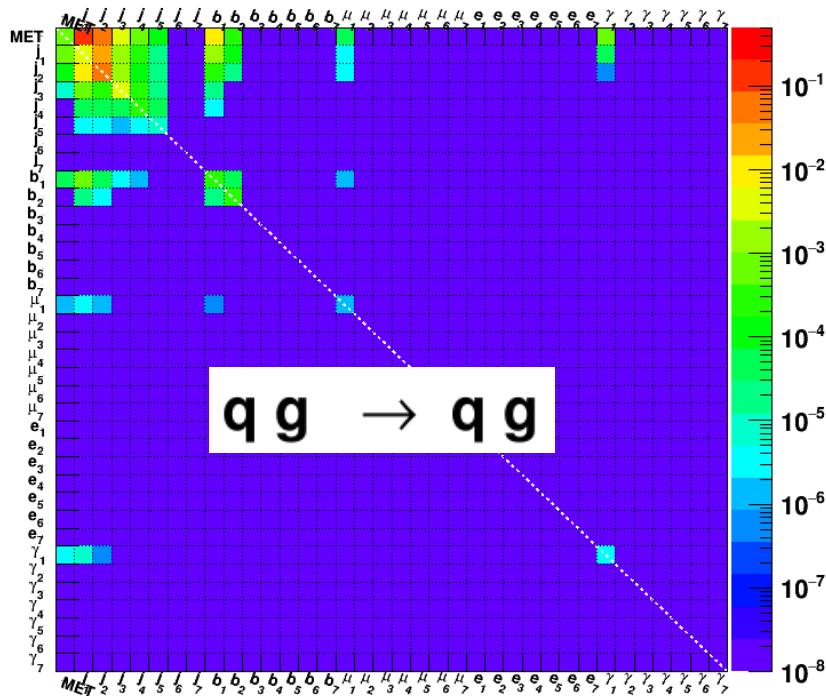
But there are many other distributions that can be used for ANN. How to choose them?



Use hand-crafted variables using Pick-and-Use approach?



Creating RMM for gg and qg events



gg process compared to qg has:

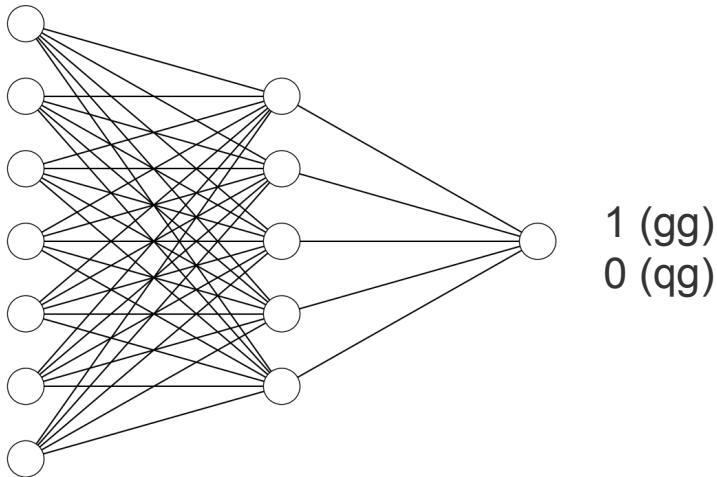
- softer p_T
- more jets
- reduced photon rate
- ..

QCD dijet separation: Challenging case!

Two approaches using machine learning:

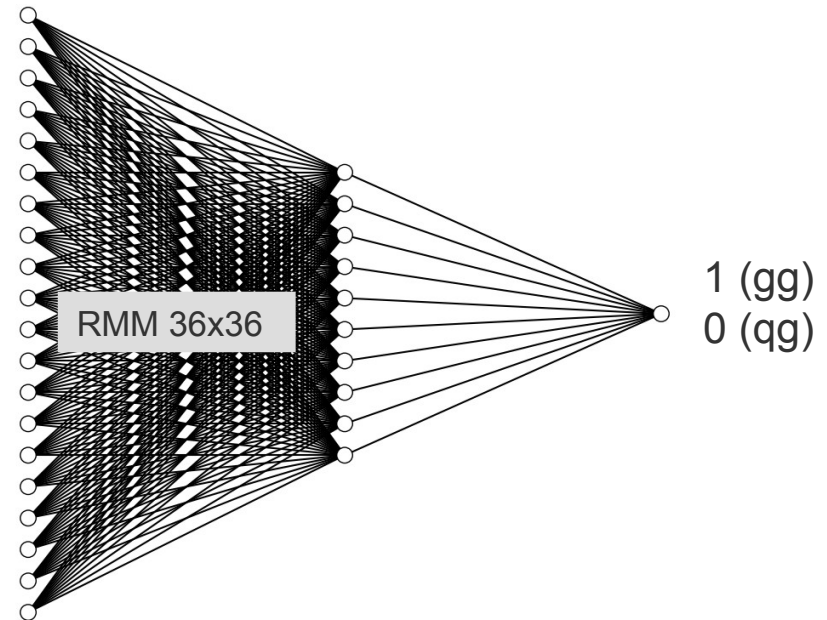
Traditional PaU

- hand-crafted input variables (7 nodes)
- hidden layer (5 nodes)
- output with 1 (gg) or 0 (qg)



RMM

- RMM matrix as input (36x36+2)
- hidden layer (200 nodes)
- output with 1 (gg) or 0 (qg)

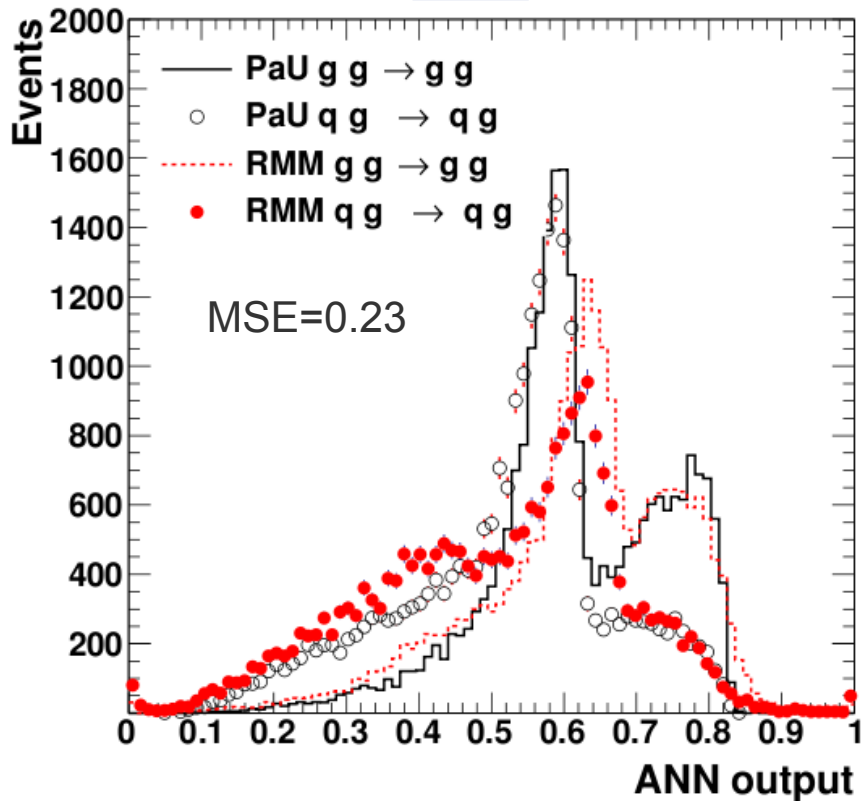


Alternatively: Use Boosted Decision tree (BDT) using PaU and RMM
100 trees, depth 7, stochastic gradient (arXiv:1609.06119)

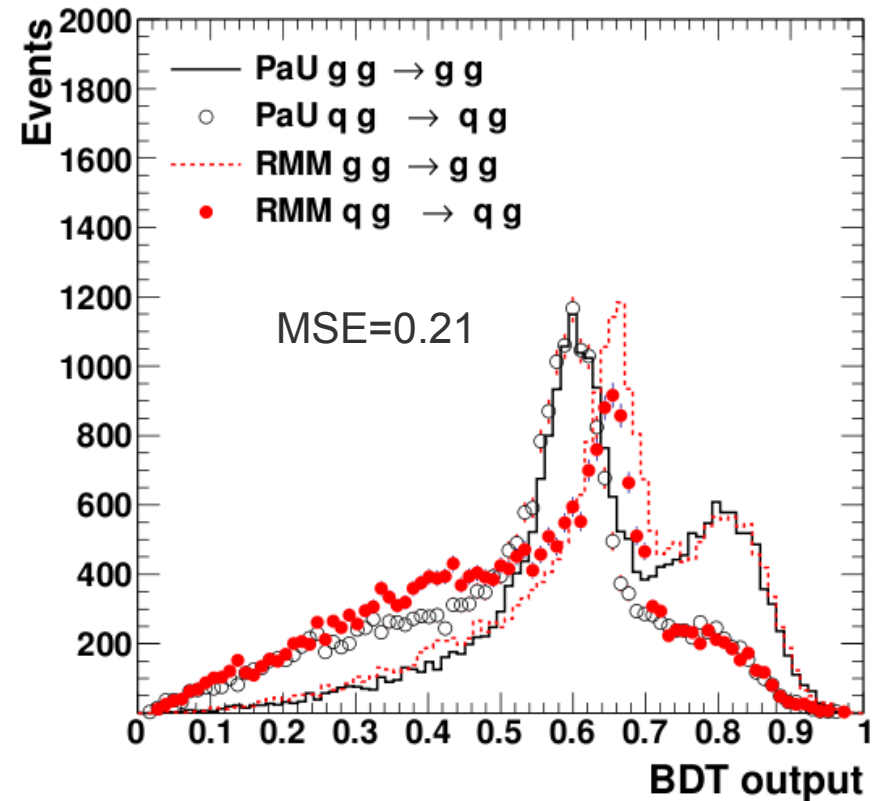


Machine learning for gg and qg separation

AAN



BDT



- Reasonable separation of gg from qg
- RMM over-performs “hand-crafted pick-and-use” (PaU) method
- RMM has separation purity 67% vs 63% for PaU assuming cut at 0.5
- BDT confirms this conclusion

Event classification using imaging of collision events. S.Chekanov (ANL) et al



Conclusions

- RMM improves event classification due comprehensive (nearly independent) single and two-particle densities for all particle types
- Same RMM transformation can be plugged into very different problems to produce good results with minimal tweaking
 - Unless you do something exotic and single and double densities of reconstructed objects are not sufficient
- Unlike hand-crafted inputs for machine learning, RMM can identify events with rather unexpected features. For example, **qg** events have an enhanced production of isolated photons. This contributes to RMM, but often escape attention
 - No need to worry about different decay channels (and their kinematics)
- If you are interested in a package that transforms events to RMM – contact me

